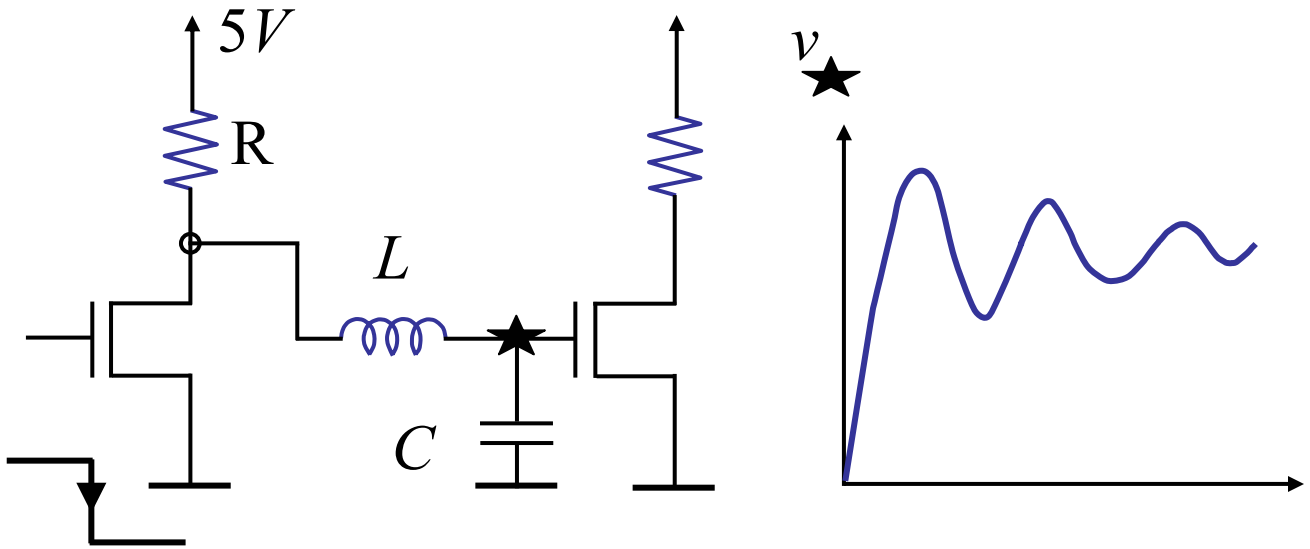


# Sinusoidal Steady State

# Review

- We now understand the why of:

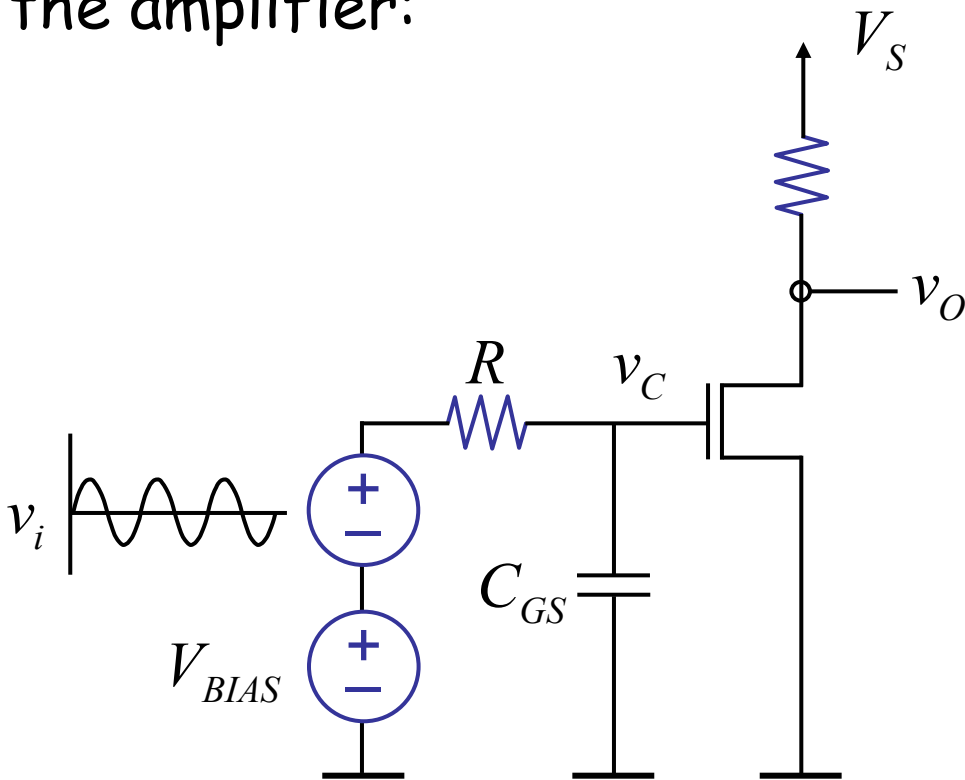


- Today, look at response of networks to sinusoidal drive.

Sinusoids important because signals can be represented as a sum of sinusoids. Response to sinusoids of various frequencies -- aka frequency response -- tells us a lot about the system

# Motivation

For motivation, consider our old friend, the amplifier:



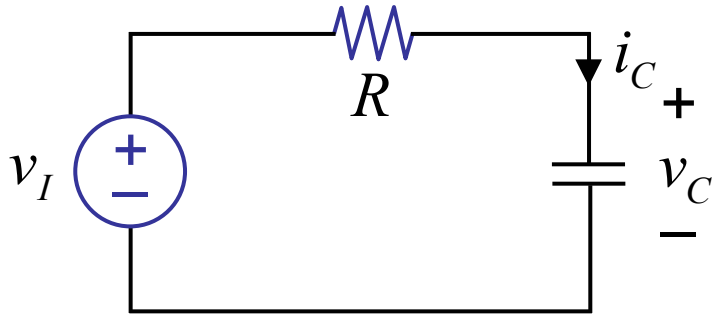
Observe  $v_o$  amplitude as the frequency of the input  $v_i$  changes. Notice it decreases with frequency.

Also observe  $v_o$  shift as frequency changes (phase).

Need to study behavior of networks for sinusoidal drive.

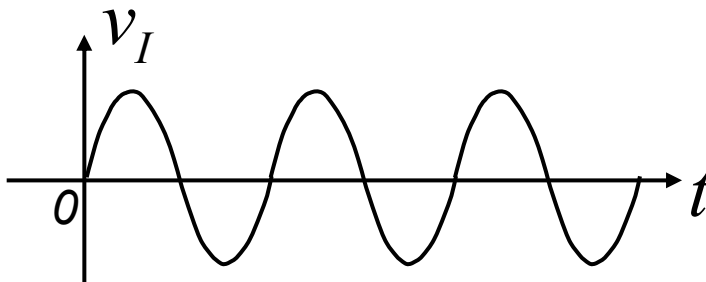
# Sinusoidal Response of RC Network

Example:



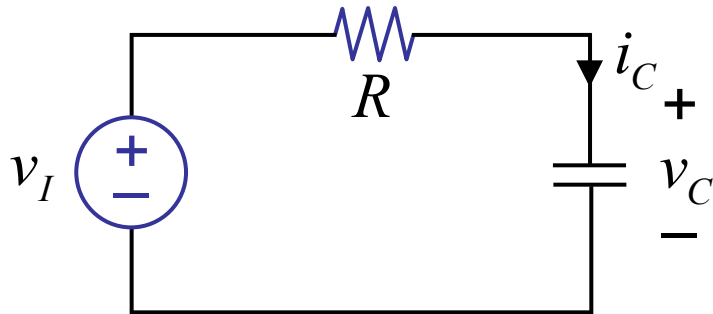
$$v_I(t) = \begin{cases} V_i \cos \omega t & \text{for } t \geq 0 \\ 0 & \text{for } t < 0 \end{cases} \quad (V_i \text{ real})$$

$$v_C(0) = 0 \quad \text{for } t = 0$$

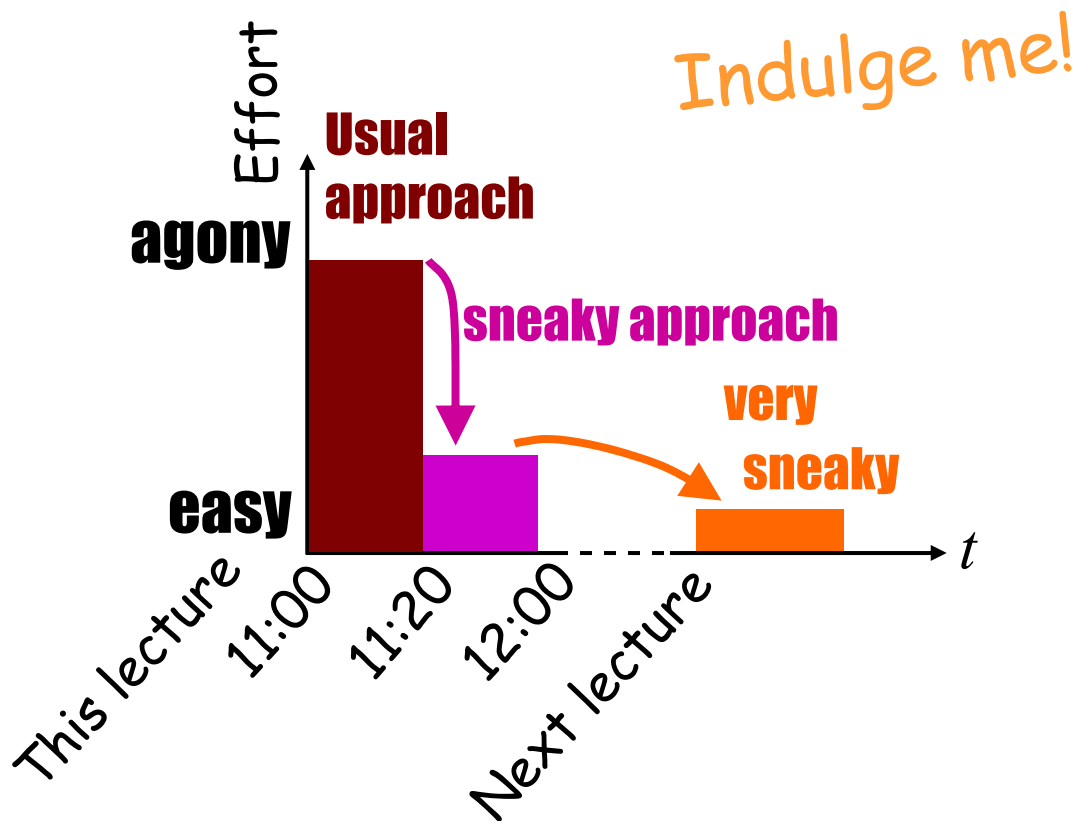


# Our Approach

Example:



Determine  $v_C(t)$



# Let's use the usual approach...

- ① Set up DE.
- ② Find  $v_P$ .
- ③ Find  $v_H$ .
- ④  $v_C = v_P + v_H$ , solve for unknowns  
using initial conditions

# Usual approach...

① Set up DE

$$\begin{aligned} RC \frac{dv_C}{dt} + v_C &= v_I \\ &= V_i \cos \omega t \end{aligned}$$

That was easy!

② Find  $v_P$

$$RC \frac{dv_P}{dt} + v_P = V_i \cos \omega t$$

First try:  $v_P = A \rightarrow$  nope

Second try:  $v_P = A \cos \omega t \rightarrow$  nope

Third try:  $v_P = A \cos(\omega t + \phi)$   
amplitude      frequency      phase

$$-RCA\omega \sin(\omega t + \phi) + A \cos(\omega t + \phi) = V_i \cos \omega t$$

$$\begin{aligned} -RCA\omega \sin \omega t \cos \phi - RCA\omega \cos \omega t \sin \phi + \\ A \cos \omega t \cos \phi - A \sin \omega t \sin \phi &= V_i \cos \omega t \end{aligned}$$

•  
• **gasp !**  
•

works, but trig nightmare!



# Let's get sneaky!

Find particular solution to another input...

$$RC \frac{dv_{PS}}{dt} + v_{PS} = v_{IS} \quad (s: \text{sneaky :-})$$
$$= V_i e^{st}$$

Try solution  $v_{PS} = V_p e^{st}$

Nice  
property  
of  
exponentials

$$RC \frac{dV_p e^{st}}{dt} + V_p e^{st} = V_i e^{st}$$

$$sRCV_p e^{st} + V_p e^{st} = V_i e^{st}$$

$$(sRC + 1)V_p = V_i$$

$$V_p = \frac{V_i}{1 + sRC}$$

$$\text{Thus, } v_{PS} = \frac{V_i}{1 + sRC} \cdot e^{st}$$

**easy!**

is particular solution to  $V_i e^{st}$

|||y  $\underbrace{\frac{V_i}{1 + j\omega RC}}_{V_p} \cdot e^{j\omega t} \rightarrow$  solution for  $V_i e^{j\omega t}$   
where we replace  $s = j\omega$

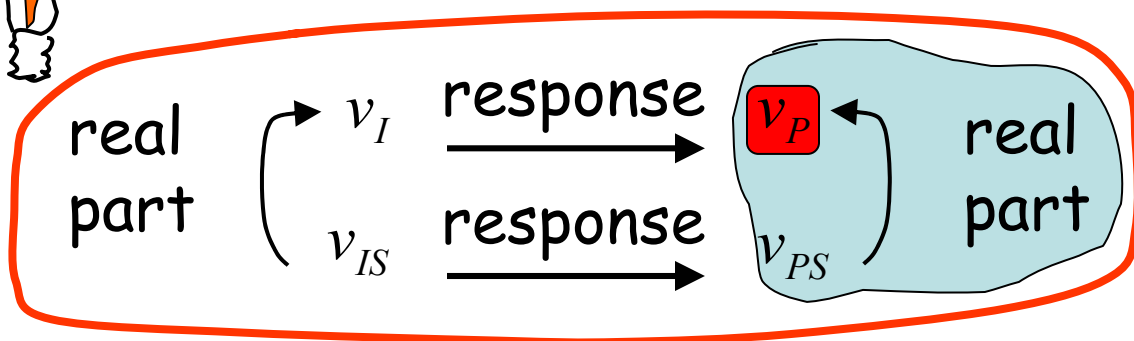
$V_p \rightarrow$  complex amplitude

② Fourth try to find  $v_P$ ..  
using the sneaky approach

**Fact 1:** Finding the response to  
 $V_i e^{j\omega t}$   
was easy.

**Fact 2:**  $v_I = V_i \cos \omega t$   
 $= \text{real}[V_i e^{j\omega t}] = \text{real}[v_{IS}]$

from Euler relation,  
 $e^{j\omega t} = \cos \omega t + j \sin \omega t$



an inverse superposition argument,  
assuming system is real, linear.

## ② Fourth try to find $v_P$ ...

so,

complex

$$v_P = \text{Re}[v_{PS}] = \text{Re}[V_p e^{j\omega t}]$$

$$= \text{Re}\left[\frac{V_i}{1 + j\omega RC} \cdot e^{j\omega t}\right]$$

$$= \text{Re}\left[\frac{V_i(1 - j\omega RC)}{1 + \omega^2 R^2 C^2} \cdot e^{j\omega t}\right]$$

$$= \text{Re}\left[\frac{V_i}{\sqrt{1 + \omega^2 R^2 C^2}} \cdot e^{j\phi} e^{j\omega t}\right], \tan \phi = -\omega RC$$

$$= \text{Re}\left[\frac{V_i}{\sqrt{1 + \omega^2 R^2 C^2}} \cdot e^{j(\omega t + \phi)}\right]$$

$$v_P = \frac{V_i}{\sqrt{1 + \omega^2 R^2 C^2}} \cdot \cos(\omega t + \phi)$$

Recall,  $v_P$  is particular response to  $V_i \cos \omega t$ .

③ Find  $v_H$

Recall,  $v_H = Ae^{\frac{-t}{RC}}$

#### ④ Find total solution

$$v_C = v_P + v_H$$

$$v_C = \frac{V_i}{\sqrt{1 + \omega^2 R^2 C^2}} \cos(\omega t + \phi) + A e^{-\frac{t}{RC}}$$

where  $\phi = \tan^{-1}(-\omega RC)$

Given  $v_C(0) = 0$  for  $t = 0$

so,

$$A = -\frac{V_i}{\sqrt{1 + \omega^2 R^2 C^2}} \cos(\phi)$$

**Done! Phew!**

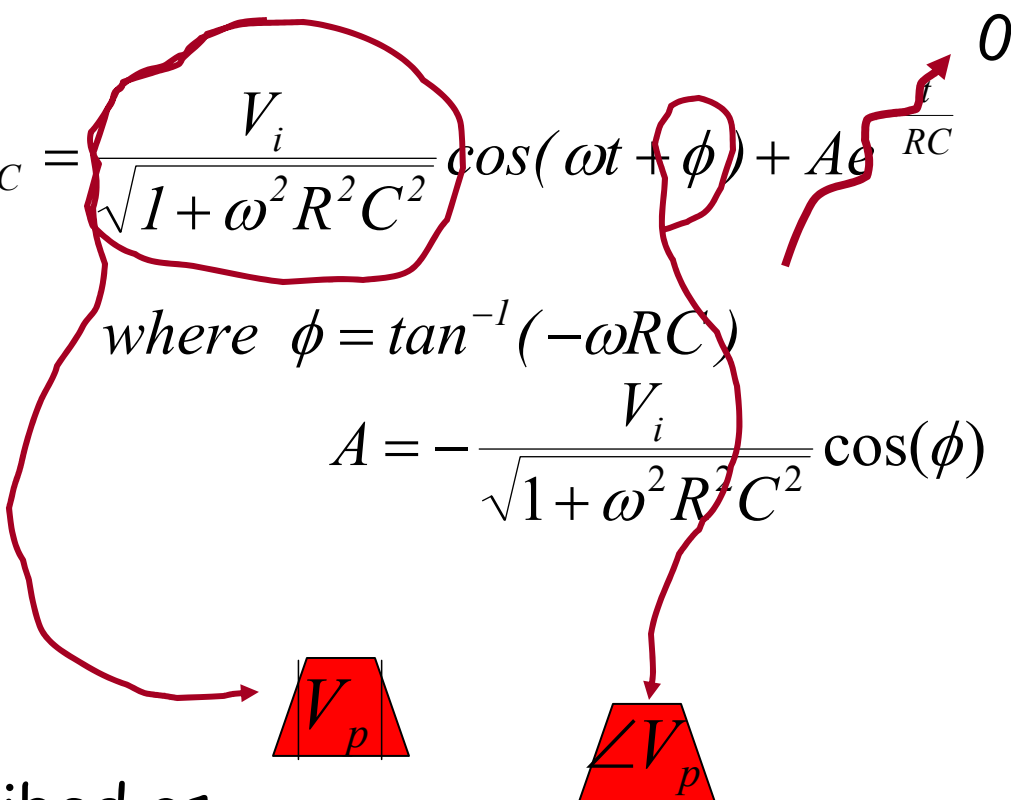
# Sinusoidal Steady State

We are usually interested only in the particular solution for sinusoids,  
i.e. *after transients have died*.

Notice when  $t \rightarrow \infty$ ,  $v_C \rightarrow v_P$  as  $e^{-\frac{t}{RC}} \rightarrow 0$

$$v_C = \frac{V_i}{\sqrt{1 + \omega^2 R^2 C^2}} \cos(\omega t + \phi) + A e^{-\frac{t}{RC}}$$


where  $\phi = \tan^{-1}(-\omega RC)$

$$A = -\frac{V_i}{\sqrt{1 + \omega^2 R^2 C^2}} \cos(\phi)$$



Described as

SSS: Sinusoidal Steady State

# Sinusoidal Steady State

All information about SSS is contained in , the complex amplitude!

Recall

$$\text{} = \frac{V_i}{1 + j\omega RC}$$

Steps ③, ④ were a waste of time!

$$\frac{V_p}{V_i} = \frac{1}{1 + j\omega RC}$$

$$\frac{V_p}{V_i} = \frac{1}{\sqrt{1 + \omega^2 R^2 C^2}} e^{j\phi} \text{ where}$$

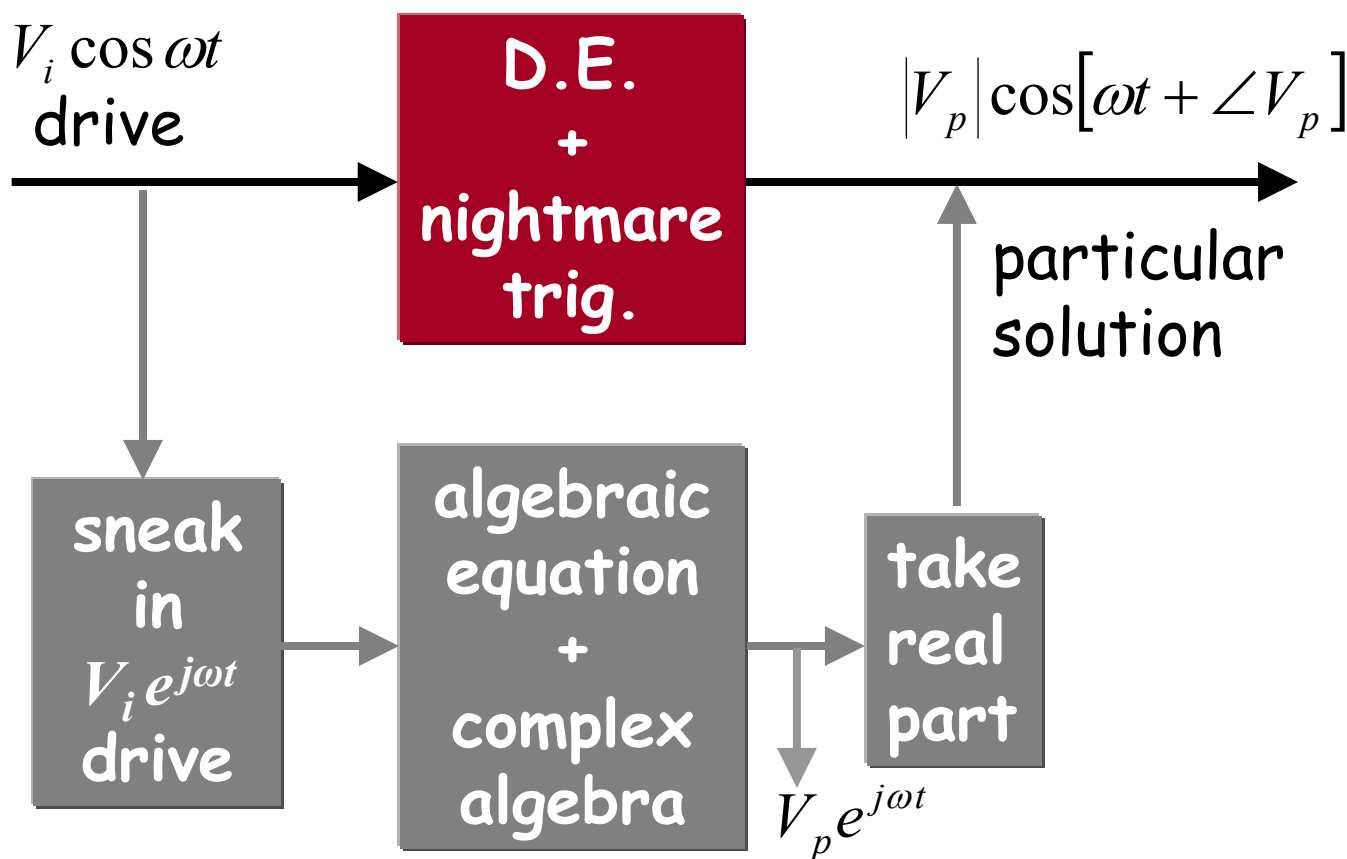
$$\phi = \tan^{-1} -\omega RC$$

magnitude  $\left| \frac{V_p}{V_i} \right| = \frac{1}{\sqrt{1 + \omega^2 R^2 C^2}}$

phase  $\phi$ :  $\angle \frac{V_p}{V_i} = -\tan^{-1} \omega RC$

# Sinusoidal Steady State

Visualizing the process of finding the particular solution  $v_P$



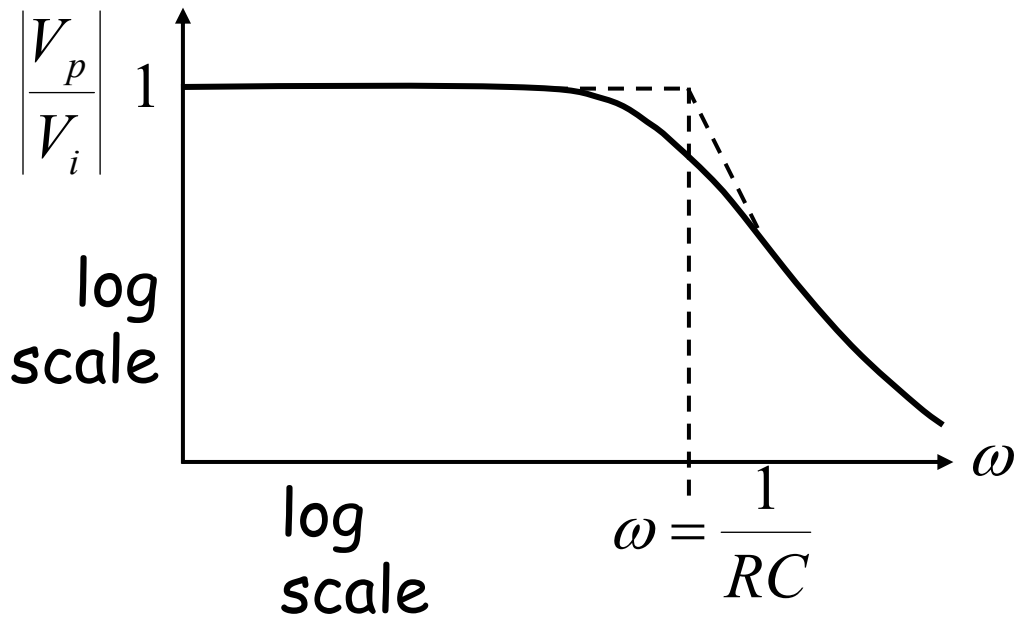
the sneaky path!



# Magnitude Plot

transfer function

$$H(j\omega) = \frac{V_p}{V_i} \qquad \left| \frac{V_p}{V_i} \right| = \frac{1}{\sqrt{1 + \omega^2 R^2 C^2}}$$



From demo: explains  $v_o$  fall off  
for high frequencies!

# Phase Plot

$$\phi = \tan^{-1} - \omega RC$$

